

# Heat equation

$$\frac{u(\vec{x}, t + \delta) - u(\vec{x}, t)}{\delta} = \Delta_h u(\vec{x})$$

$$u(\vec{x}, 0) = u_0(\vec{x})$$

$$u(\vec{x}, t + \delta) = u(\vec{x}, t) + \delta \Delta_h u(\vec{x}, t)$$

# Wave equation

$$\frac{u(\vec{x}, t + \delta) - 2u(\vec{x}, t) + u(\vec{x}, t - \delta)}{\delta^2} = \Delta_h u(\vec{x}, t)$$

$$\frac{\partial u}{\partial t}(\vec{x}, 0) = 0$$

$$u(\vec{x}, 0) = u_0(\vec{x}), \text{ and } u(\vec{x}, \delta) = u_0(\vec{x})$$

$$u(\vec{x}, t + \delta) = 2u(\vec{x}, t) - u(\vec{x}, t - \delta) + \delta^2 \Delta_h u(\vec{x}, t)$$

# Matrix representation in 1-D

- 1-D Poisson

$$\frac{1}{h^2} \begin{pmatrix} -2 & 1 & & & & & \\ 1 & -2 & 1 & & & & \\ & 1 & -2 & 1 & & & \\ & & \dots & \dots & \dots & & \\ & & & 1 & -2 & 1 & \\ & & & & 1 & -2 & \end{pmatrix}$$

With  $u$  represented as a vector and  $h^2\Delta_h$  as a matrix  $A$ , the Poisson problem becomes

$$Au = b$$

where  $b$  is a vector (the same size as  $u$ ) containing the values of  $h^2f(x)$  at the interior mesh points. The first and last components of  $b$  would also include any nonzero boundary values.

$$u = A \setminus b$$

In Matlab this takes advantage of the sparse structure of  $A$

# Matrix representation in 2D

- Need to map 2D domain in to 1-D element
- Assume a L shaped region

L =

0	0	0	0	0	0	0	0	0	0	0
0	1	5	9	13	17	21	30	39	48	0
0	2	6	10	14	18	22	31	40	49	0
0	3	7	11	15	19	23	32	41	50	0
0	4	8	12	16	20	24	33	42	51	0
0	0	0	0	0	0	25	34	43	52	0
0	0	0	0	0	0	26	35	44	53	0
0	0	0	0	0	0	27	36	45	54	0
0	0	0	0	0	0	28	37	46	55	0
0	0	0	0	0	0	29	38	47	56	0
0	0	0	0	0	0	0	0	0	0	0

- Boundary conditions on 1-4, 5,8,9,12,13,16, 17, 20, 21, 29,30,38,39,47, 48-56

- Write the operator as  $Au=f$
- Consider node 43

33	42	51
34	43	52
35	44	53

$$h^2 \Delta_h u(43) = u(34) + u(42) + u(44) + u(52) - 4u(43)$$

- So the elements of the matrix  $a$  for the row corresponding to 43 are

$$a_{43,34} = a_{43,42} = a_{43,44} = a_{43,52} = 1, \text{ and } a_{43,43} = -4$$

- All other entries are zero
- Each row for an internal node can have at most 5 non-zero entries

# Stability for time-dependent equations

- Can study stability as we did for ODEs
- Gives a condition relating timestep size and mesh size
- Heat equation

$$u^{(k+1)} = u^{(k)} + \sigma A u^{(k)}$$

$$u^{(k+1)} = M u^{(k)}$$

$$\sigma = \frac{\delta}{h^2}$$

$$M = I + \sigma A$$

- For stability we require in one dimension
- So mesh size and time step cannot be chosen independently

$$\sigma \leq \frac{1}{2}$$

In two dimensions

$$\sigma \leq \frac{1}{4}$$

# Wave equation stability criteria

$$u^{(k+1)} = 2u^{(k)} - u^{(k-1)} + \sigma Au^{(k)}$$

$$\sigma = \frac{\delta^2}{h^2}$$

For the wave equation in one dimension

$$\sigma \leq 1$$

In two dimensions

$$\sigma \leq \frac{1}{2}$$

# Wave equation

- Separation of variables

$$\frac{\partial^2 u}{\partial t^2} = \Delta u$$

- Gives the “Helmholtz” equation for  $v$

$$u(\vec{x}, t) = \cos(\sqrt{\lambda} t) v(\vec{x})$$

- Solving for  $v$

$$\Delta v + \lambda v = 0$$

- So overall solutions is

In one dimension

$$v_k(x) = \sin(kx)$$

$$v_k(\pi) = 0$$

$k$  must be an integer

$$\lambda_k = k^2.$$

$$\begin{aligned} u(x, t) &= \sum_k a_k \cos(kt) \sin(kx) \\ &= \sum_k a_k \cos(\sqrt{\lambda_k} t) v_k(x) \end{aligned}$$



